

- BIBO stabilitás (nagyarányú rövidek) } $v|z+K$
 → LTI stabilitás (nagyarányú rövidek) } Csisz + Neuton pld.



$$\dot{x} = \begin{pmatrix} -1 & 0 \\ 1 & 1 \end{pmatrix} x + \begin{pmatrix} 1 \\ 0 \end{pmatrix} u$$

$$y = (1 \ 0) x$$

nem stabil

$$H(s) = \frac{1}{s+1}$$

BIBO stab

→ Zoltka-Volterra:

$$\dot{x}_1 = +7x_1 - 4x_1^2 - 3x_1x_2 \quad (\text{nyúlók / szélvendég})$$

$$\dot{x}_2 = -3x_2 + 2x_1x_2 + x_2^2 \quad (\text{farkas / ragadozó})$$

$$\hookrightarrow x_{(0)}^* = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \rightarrow \text{instab}$$

$$\hookrightarrow x_{(1)}^* = \begin{pmatrix} 1 \\ 1 \end{pmatrix} \rightarrow \text{stabil}$$

→ Joga (síma) + Zjopunov függvény

$$\theta = \omega$$

$$\dot{\omega} = -\frac{g}{l} \sin \theta - b\omega$$

$$V = mgl(1 - \cos \theta) + \frac{m l^2 \omega^2}{2}$$

(Szemléletes + vides)

→ (Joga Hamilton rendszer) ← erre talán nem lenn idek

$$\dot{q} = \frac{\partial H}{\partial p}$$

$$\dot{p} = -\frac{\partial H}{\partial q} - b p$$

$$\text{ahol } p = \frac{\partial \mathcal{L}}{\partial \dot{q}}$$

→ Van der Pol oscillator (paramétrus Zj.p.) + optium

$$\dot{x}_1 = -x_2$$

$$\dot{x}_2 = +x_1 + (1 - x_1^2)x_2$$

$$V = x^T \begin{pmatrix} a & b \end{pmatrix} x$$

Zoltea-Volterra: $x = \text{diag}(x)(Ax + b) \Rightarrow x^* = -A^{-1}b$
 $b = -Ax^*$

$$A = \begin{pmatrix} -2 & -3 \\ 1.4 & 1 \end{pmatrix} \Rightarrow b = -\begin{pmatrix} -2 & -3 \\ 1.4 & 1 \end{pmatrix} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = -\begin{pmatrix} -5 \\ 2.4 \end{pmatrix}$$

$$x^* = \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

$$\begin{cases} \dot{x}_1 = x_1(-2x_1 - 3x_2 + 5) \\ \dot{x}_2 = x_2(1.4x_1 + x_2 - 2.4) \end{cases}$$

nyugale
ragadozó

Egyensúlyi pontok:

$$x_{(0)}^* = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \quad x_{(1)}^* = \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

átírva:

$$\begin{cases} \dot{x}_1 = -2x_1^2 - 3x_1x_2 + 5x_1 \\ \dot{x}_2 = 1.4x_1x_2 + x_2^2 - 2.4x_2 \end{cases}$$

$$f(x) = \begin{pmatrix} -2x_1^2 - 3x_1x_2 + 5x_1 \\ 1.4x_1x_2 + x_2^2 - 2.4x_2 \end{pmatrix}$$

$$\frac{\partial f}{\partial x} = \begin{pmatrix} -4x_1 - 3x_2 + 5 & -3x_1 \\ 1.4x_2 & 1.4x_1 + 2x_2 - 2.4 \end{pmatrix}$$

$$\frac{\partial f}{\partial x}(0,0) = \begin{pmatrix} 5 & 0 \\ 0 & -2.4 \end{pmatrix}$$

tehát az $x_{(0)}^* = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$ nem stabil egy. pont

$$\frac{\partial f}{\partial x}(1,1) = \begin{pmatrix} -2 & -3 \\ 1.4 & 1 \end{pmatrix}$$

$$\begin{aligned} \det \begin{pmatrix} s+2 & 3 \\ -1.4 & s-1 \end{pmatrix} &= s^2 + s - 2 + 4.2 \\ &= s^2 + s + 2.2 \end{aligned}$$

$$s_{1,2} = -\frac{1}{2} \pm 1.396i$$

stabil

$$\Delta = 1 - 8.8 = -7.8$$

Megjegyzés: $A = \begin{pmatrix} -4 & -3 \\ 2 & 1 \end{pmatrix}$ -re $s_{1,2} \in \{-1, -2\}$.

2020b

Nonlineáris rendszer lokális stabilitása.

$$\ddot{\theta} + b\dot{\theta} + \frac{g}{l} \sin \theta = 0$$

Pendulum ← Lyapunov fu.

$$\begin{cases} x_1 = \theta \\ x_2 = \dot{\theta} \end{cases} \Rightarrow \begin{cases} \dot{x}_1 = \dot{\theta} = x_2 \\ \dot{x}_2 = \ddot{\theta} = -\frac{g}{l} \sin \theta - b\dot{\theta} = -\frac{g}{l} \sin x_1 - b x_2 \end{cases}$$

$$x = \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} \text{ allora } \dot{x} = f(x), \text{ ahol } f(x) = \begin{pmatrix} x_2 \\ -\frac{g}{l} \sin x_1 - b x_2 \end{pmatrix}$$

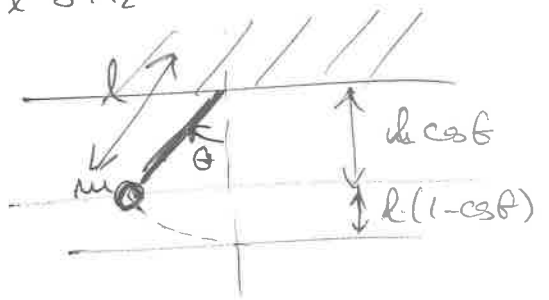
$$\text{Leggiamo } V(x_1, x_2) = \frac{m l^2 x_2^2}{2} + m g l (1 - \cos x_1)$$

$$\dot{V} = \frac{\partial V}{\partial x_1} \dot{x}_1 + \frac{\partial V}{\partial x_2} \dot{x}_2 = \left(\frac{\partial V}{\partial x} f(x) \right)$$

$$= (m g l \sin x_1) \cdot x_2 + (m l^2 x_2) \cdot \left(-\frac{g}{l} \sin x_1 - b x_2 \right)$$

$$= m g l x_2 \sin x_1 - m g l x_2 \sin x_1 - m l^2 b x_2^2$$

$$= -m l^2 b x_2^2$$



$$\dot{V} = -m l^2 b x_2^2 < 0$$

damping coefficient!

The Hamiltonian mechanics can ask why the total energy is the best

Lyapunov func?

$$\begin{cases} \dot{q} = \frac{\partial \mathcal{H}}{\partial p} \\ \dot{p} = -\frac{\partial \mathcal{H}}{\partial q} - b p \end{cases} \quad (\text{Hamiltonian model of the pendulum})$$

$$\text{Let } V = \mathcal{H} = \frac{p^2}{2m l^2} + m g l (1 - \cos q)$$

$$\dot{V} = \frac{\partial V}{\partial p} \dot{p} + \frac{\partial V}{\partial q} \dot{q} = \frac{\partial \mathcal{H}}{\partial p} \dot{p} + \frac{\partial \mathcal{H}}{\partial q} \dot{q} =$$

$$= \frac{\partial \mathcal{H}}{\partial p} \left(-\frac{\partial \mathcal{H}}{\partial q} - b p \right) + \frac{\partial \mathcal{H}}{\partial q} \frac{\partial \mathcal{H}}{\partial p} = -m l^2 b p^2 < 0$$

miss each other ↙

$$Q = T - V = \frac{ml^2\omega^2}{2} - mgl(1 - \cos\theta) = \frac{ml^2\dot{q}^2}{2} - mgl(1 - \cos q)$$

leggjum $q = \theta$ (jafn: allt. koordinata) } ↑

leggjum $\dot{q} = \dot{\theta} = \omega$ (allt. sebessög) } ↑

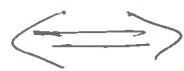
leggjum $p = \frac{\partial Q}{\partial \dot{q}} = ml^2\dot{q}$ (alteldingasettitt mamentum)

$$\begin{aligned} \text{leggjum } \mathcal{H} &= \dot{q} \cdot p - Q = \dot{q} p - \left(\frac{ml^2\dot{q} \cdot \dot{q}}{2} - mgl(1 - \cos q) \right) \\ &= \frac{ml^2\dot{q} \cdot \dot{q}}{2} + mgl(1 - \cos q) \\ &= \frac{p^2}{2ml^2} + mgl(1 - \cos q) \end{aligned}$$

↑ Sum of kinetic and potential energy
in the terms of the generalised coordinate (q)
→ momentum (p)

First order DE of the motion:

$$\begin{cases} \dot{q} = + \frac{\partial \mathcal{H}}{\partial p} \\ \dot{p} = - \frac{\partial \mathcal{H}}{\partial q} \end{cases}$$



Euler-Lagrange equations:

$$\frac{d}{dt} \frac{\partial Q}{\partial \dot{q}} - \frac{\partial Q}{\partial q} = 0$$



$$\frac{d}{dt} (ml^2\dot{q}) + mgl \sin q = 0$$



$$ml^2\ddot{q} + mgl \sin q = 0$$



$$\ddot{q} + \frac{g}{l} \sin q = 0$$

$$\begin{cases} \dot{q} = \frac{p}{ml^2} \\ \dot{p} = -mgl \sin q \end{cases}$$

$$\ddot{q} = \frac{\dot{p}}{ml^2} = -\frac{g}{l} \sin q$$

$$\ddot{q} + \frac{g}{l} \sin q = 0$$

20206

Hamiltonian model
of pendulum.

Von der Pol Ljapunov's Lyapunov fu. (2020b)

$$\begin{cases} \dot{x}_1 = -x_2 \\ \dot{x}_2 = x_1 - (1-x_1^2)x_2 \end{cases} \Rightarrow \begin{cases} \dot{x}_1 = -x_2 \\ \dot{x}_2 = x_1 - x_2 + x_1^2 x_2 \end{cases}$$

Leggen $V(x) = x^T \begin{pmatrix} a & b \\ 0 & 0 \end{pmatrix} x$ parametrisiertes Lyapunov fu. gelöst
 $a, b = ?$ n.h. (local abh. lokal stabil) $V(x)$ L.f. legen

$$\begin{aligned} \frac{\partial V}{\partial x} f(x) &= x^T \begin{pmatrix} a & b \\ 0 & 0 \end{pmatrix} \begin{pmatrix} -x_2 \\ x_1 - x_2 + x_1^2 x_2 \end{pmatrix} = \\ &= (x_1 a \quad x_2 b) \cdot \begin{pmatrix} \dots \\ \dots \end{pmatrix} = \underbrace{-x_1 x_2 a + x_1 x_2 b - x_2^2 b + x_1^2 x_2^2 b}_{\text{da } a=b \text{ alle ersele}} \end{aligned}$$

Kürztheile

Leggen $a = b = 1$, also $V(x) = x^T \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} x = x^T x = x_1^2 + x_2^2$

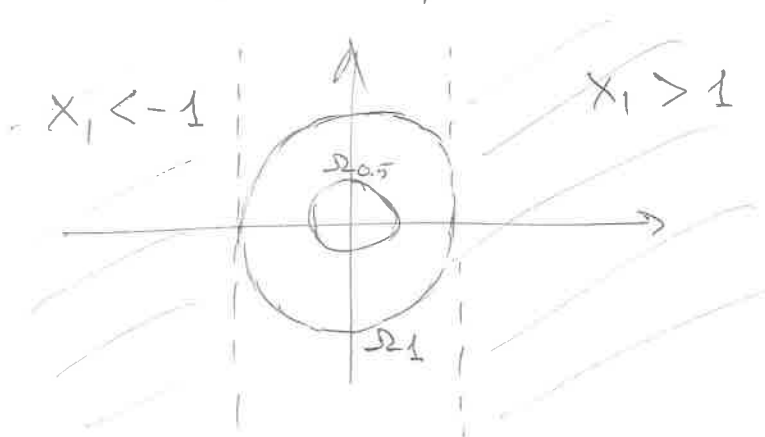
etlich:

$$\frac{\partial V}{\partial x} f(x) = -x_2^2 + x_1^2 x_2^2 = -x_2^2 (1 - x_1^2) < 0 \quad \text{da}$$

gilt $|x| < 1$ ←

Invarianz forstreuung (stabil region):

$$\Omega_c = \{x \in \mathbb{R}^2 \mid V(x) \leq c\} \subseteq \{x \in \mathbb{R}^2 \mid |x| < 1\}$$



↳ Menge der stabil region
 (also $\frac{\partial V}{\partial x} f(x) < 0$)

$\Omega_{1-\epsilon}$ ist jedes gegebenes benue von

2020b.

Von der Pol
 Ljapunov's Lyapunov fu.