

Computer controlled systems

Crane model

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1 Crane model (rakodó darumodell)

Consider the following mathematical model of a crane machine. For the sake of simplicity, we restrict the motion of the carried weight to the (y, z) plane. In the model we identify the following time-dependent (state-) variables:

- $R = R(t)$ denotes the actual position of the car on the rail,
- $L = L(t)$ denotes the actual length of the wire,
- $\theta = \theta(t)$ denotes the actual angle of the wire with the vertical.

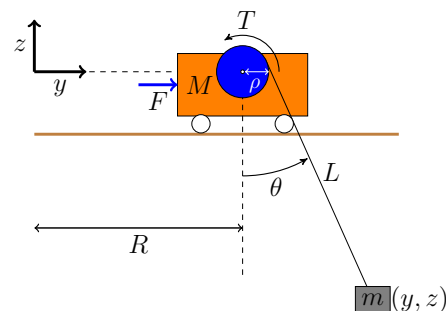


Figure 1. Kinetic model of the crane machine.

Known parameters (constants):

- ρ is the radius of the pulley¹. In the geometrical expression the radius of the pulley is neglected, e.g. if the angle is $\theta = 0$ and the position of the car is $R = 0$ then the position of the weight along the y axis is considered to be R .
- M is the mass of the car
- m is the mass of the lifted weight
- J is the moment of inertia² of the pulley.

Manipulated inputs are the following:

- F is the driving force applied to the car
- T is the torque applied to the pulley

Measured quantities: $R(t)$ and $L(t)$.

Tekintse az ábrán látható rakodó darut. Az egyszerűség kedvéért a teher mozgását az (y, z) függőleges síkra korlátozzuk. Jelölje R a sínen mozgó kocsi y irányú pozícióját, L a sodrony hosszát és θ a sodrony függőlegessel bezárt szögét. Ismert geometriai paraméter a sodronyt tekerceselő dob sugara, melyet ρ -val jelöltünk. (A geometriai összefüggésekben a dob sugarát elhanyagoljuk, azaz $\theta = 0$ esetben $R = y$). Szintén ismertek az inercia paraméterek: M a kocsi tömege; m a mozgatott teher tömege és J a sodronyt tekerceselő hajtás és a dob tehetetlenségi nyomatéka.

A beavatkozó jelek:

- F a kocsira ható erő
- T a sodronyt tekerceselő dombra ható forgatónyomaték

A mért kimeneti változók: R és L .

The system's kinetic energy is a composition of the followings:

¹sheave or drum: tekerceselő csiga vagy tekerceselő dob

²tehetetlenségi nyomaték

1. the kinetic energy of M is $T_M = \frac{M\dot{R}^2}{2}$
2. the kinetic energy of m is $T_m = \frac{mv^2}{2}$
3. the kinetic energy of the pulley is $T_J = \frac{J\dot{\theta}^2}{2} = \frac{J\dot{L}^2}{2\rho^2}$

The system's potential energy is the potential energy of m , that is $V_m = mgL(1 - \cos \theta)$.

The system's Lagrangian is

$$\mathcal{L} = T - V = \frac{M\dot{R}^2}{2} + \frac{mv^2}{2} + \frac{J\dot{L}^2}{2\rho^2} - mgL(1 - \cos \theta) \quad (1)$$

The velocity \mathbf{v} of m has the following components (see Figure 1):

$$\mathbf{v} = L\dot{\theta}\mathbf{e}_t + \dot{R}\mathbf{e}_x + \dot{L}\mathbf{e}_n \quad (2)$$

Since $\mathbf{e}_t \perp \mathbf{e}_n$, the square of the norm of \mathbf{v} can be computed in the following way

$$\begin{aligned} v = \|\mathbf{v}\|^2 &= \langle \mathbf{v}, \mathbf{v} \rangle = L^2\dot{\theta}^2 + \dot{R}^2 + \dot{L}^2 + 2L\dot{\theta}\dot{R}\langle \mathbf{e}_t, \mathbf{e}_x \rangle + 2\dot{R}\dot{L}\langle \mathbf{e}_n, \mathbf{e}_x \rangle \\ &= L^2\dot{\theta}^2 + \dot{R}^2 + \dot{L}^2 + 2L\dot{\theta}\dot{R}\cos \theta + 2\dot{R}\dot{L}\cos \left(\frac{\pi}{2} - \theta \right) \\ &= L^2\dot{\theta}^2 + \dot{R}^2 + \dot{L}^2 + 2L\dot{\theta}\dot{R}\cos \theta + 2\dot{R}\dot{L}\sin \theta \end{aligned} \quad (3)$$

Therefore, the Lagrangian can be written in the following form:

$$\mathcal{L} = \frac{(M+m)\dot{R}^2}{2} + \frac{mL^2\dot{\theta}^2}{2} + \frac{m\dot{L}^2}{2} + mL\dot{\theta}\dot{R}\cos \theta + m\dot{R}\dot{L}\sin \theta + \frac{J\dot{L}^2}{2\rho^2} - mgL(1 - \cos \theta) \quad (4)$$

The Euler-Lagrange equations are the following:

$$\begin{aligned} 0 &= \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{R}} - \frac{\partial \mathcal{L}}{\partial R} = \frac{d}{dt} \left((M+m)\dot{R} + mL\dot{\theta}\cos \theta + m\dot{L}\sin \theta \right) \\ &= (M+m)\ddot{R} + m\dot{L}\dot{\theta}\cos \theta + mL\ddot{\theta}\cos \theta - mL\dot{\theta}^2\sin \theta + m\ddot{L}\sin \theta + m\dot{L}\dot{\theta}\cos \theta \\ &= (M+m)\ddot{R} + 2m\dot{L}\dot{\theta}\cos \theta + mL\ddot{\theta}\cos \theta - mL\dot{\theta}^2\sin \theta + m\ddot{L}\sin \theta \end{aligned} \quad (A1)$$

Second equation:

$$\begin{aligned} 0 &= \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{L}} - \frac{\partial \mathcal{L}}{\partial L} = \frac{d}{dt} \left(m\dot{L} + m\dot{R}\sin \theta + \frac{J\dot{L}}{\rho^2} \right) - mL\dot{\theta}^2 - m\dot{\theta}\dot{R}\cos \theta + mg(1 - \cos \vartheta) \\ &= \left(m + \frac{J}{\rho^2} \right) \ddot{L} + m\ddot{R}\sin \theta + \underbrace{m\dot{R}\dot{\theta}\cos \theta}_{\text{red}} - mL\dot{\theta}^2 - \underbrace{m\dot{\theta}\dot{R}\cos \theta}_{\text{red}} + mg(1 - \cos \vartheta) \\ &= \left(m + \frac{J}{\rho^2} \right) \ddot{L} + m\ddot{R}\sin \theta - mL\dot{\theta}^2 + mg(1 - \cos \vartheta) \end{aligned} \quad (A2)$$

Third equation:

$$\begin{aligned} 0 &= \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\theta}} - \frac{\partial \mathcal{L}}{\partial \theta} = \frac{d}{dt} \left(mL^2\dot{\theta} + mL\dot{R}\cos \theta \right) + mL\dot{\theta}\dot{R}\sin \theta - m\dot{R}\dot{L}\cos \theta + mgL\sin \theta \\ &= 2mL\dot{L}\dot{\theta} + mL^2\ddot{\theta} + \underbrace{m\dot{L}\dot{R}\cos \theta}_{\text{blue}} + mL\ddot{R}\cos \theta - \underbrace{mL\dot{R}\dot{\theta}\sin \theta}_{\text{red}} + \underbrace{mL\dot{\theta}\dot{R}\sin \theta}_{\text{red}} - \underbrace{m\dot{R}\dot{L}\cos \theta}_{\text{blue}} + mgL\sin \theta \\ &= 2mL\dot{L}\dot{\theta} + mL^2\ddot{\theta} + mL\ddot{R}\cos \theta + mgL\sin \theta \end{aligned}$$

Dividing by L we get:

$$0 = 2m\dot{L}\dot{\theta} + mL\ddot{\theta} + m\ddot{R}\cos \theta + mg\sin \theta \quad (A3)$$

In equilibrium point the torque is $T_0 = mg\rho$. Let $T = T_0 + \tau$ be the net torque. Furthermore, we consider

Since $F(x_0, u_0) = 0$, the linearized model is $\dot{x} = Ax + Bu$, where:

$$A = \left[\frac{\partial F(x, u)}{\partial x} \right]_{\substack{x=0 \\ u=0}} = \left(\begin{array}{cc|ccc} 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & \frac{gm}{M} & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & -\frac{g(M+m)}{L_0M} & 0 \end{array} \right) \quad B = \left[\frac{\partial F(x, u)}{\partial u} \right]_{\substack{x=0 \\ u=0}} = \left(\begin{array}{c|c} 0 & 0 \\ \hline 0 & -\frac{\rho}{m\rho^2+J} \\ \hline 0 & 0 \\ \frac{1}{M} & 0 \\ 0 & 0 \\ -\frac{1}{L_0M} & 0 \end{array} \right) \quad (16)$$

As one can immediately observe, we obtained two decoupled subsystems in the linearized model:

$$\begin{aligned} \dot{\xi}_1 &= A_1 \xi_1 + B_1 \tau, \quad \text{where} \quad A_1 = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}, \quad B_1 = \begin{pmatrix} 0 \\ -\frac{\rho}{m\rho^2+J} \end{pmatrix} \\ \dot{\xi}_2 &= A_2 \xi_2 + B_2 f, \quad \text{where} \quad A_2 = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & \frac{gm}{M} & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & -\frac{g(M+m)}{L_0M} & 0 \end{pmatrix}, \quad B_2 = \begin{pmatrix} 0 \\ \frac{1}{M} \\ 0 \\ -\frac{1}{L_0M} \end{pmatrix} \end{aligned} \quad (17)$$

The new state vectors ξ_1 and ξ_2 are the following:

$$\xi_1 = \begin{pmatrix} l \\ \dot{l} \end{pmatrix} = \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}, \quad \xi_2 = \begin{pmatrix} r \\ \dot{r} \\ \theta \\ \dot{\theta} \end{pmatrix} = \begin{pmatrix} x_3 \\ x_4 \\ x_5 \\ x_6 \end{pmatrix} \quad (18)$$